

# PARTON DISTRIBUTIONS AND $\Lambda_{\text{QCD}}$ AT HERA

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# 1. INTRODUCTION

$$\text{HERA: } e^{\pm} p \rightarrow \begin{pmatrix} \bar{\nu}_e \\ e^{\pm} \end{pmatrix} X$$

$$S = 4 \cdot 30 \cdot 820 \text{ GeV}^2$$

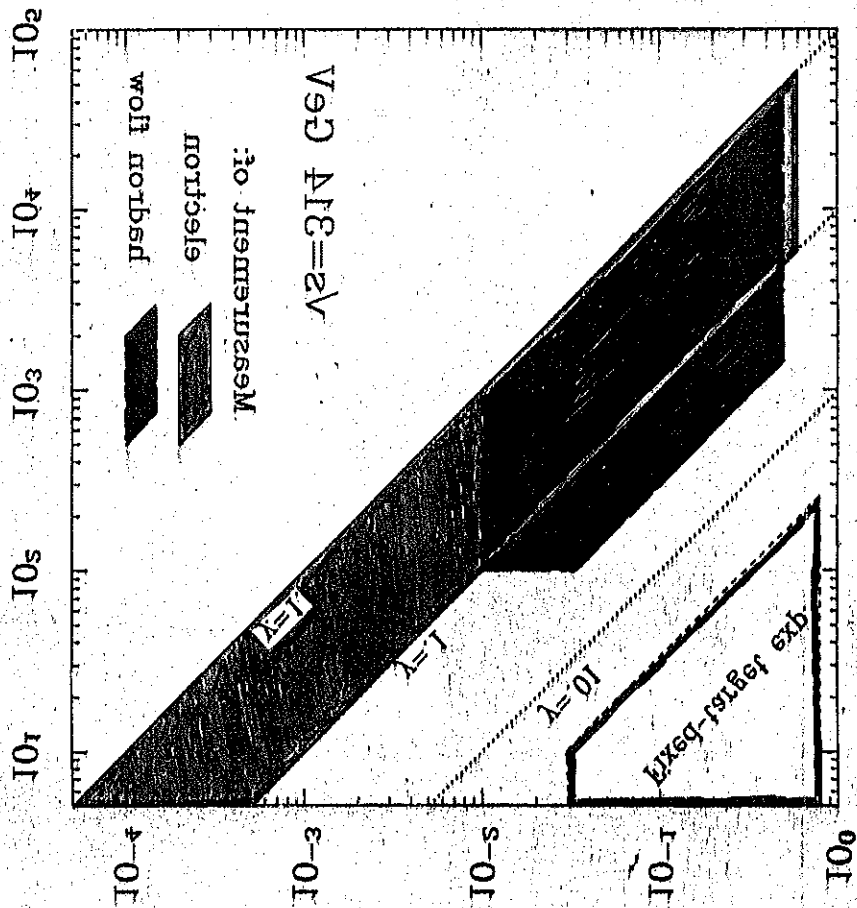
$$\mathcal{L} = 1.5 \cdot 10^{31} \text{ cm}^{-2} \text{ sec}^{-1}$$

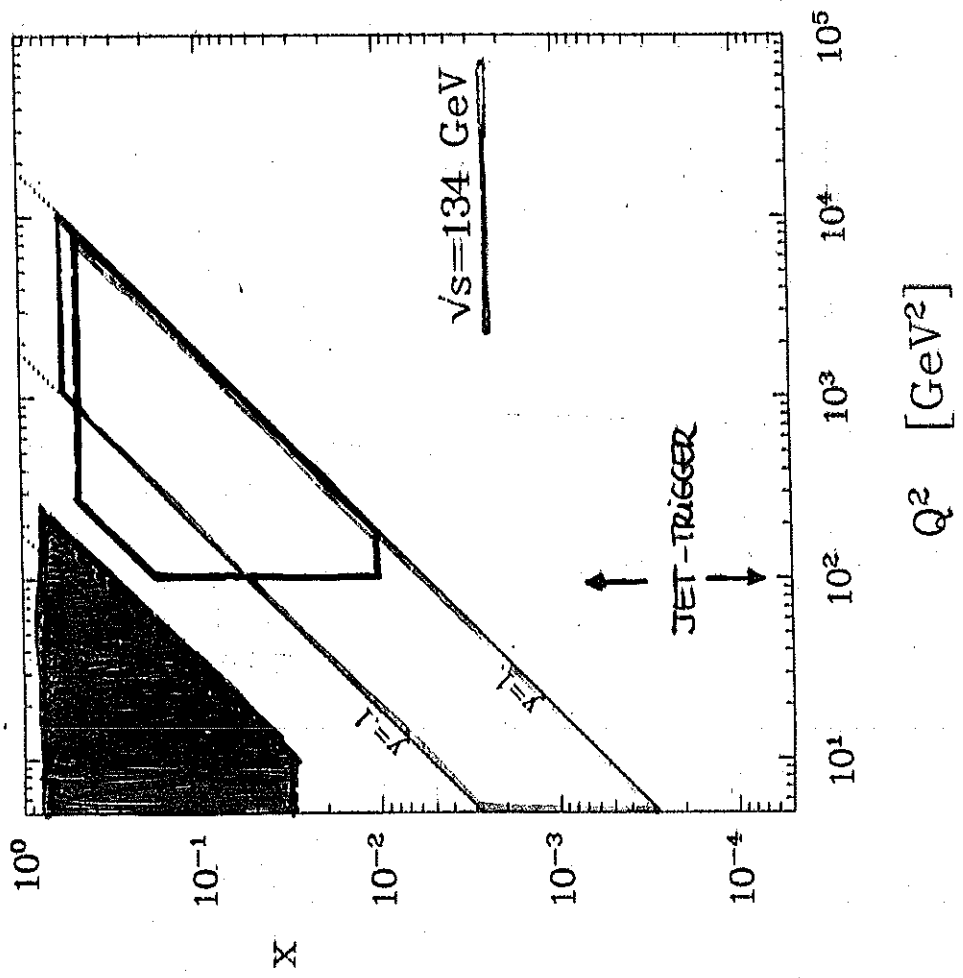
$$C \rightarrow 200 \text{ pb}^{-1} \approx 150 \text{ fully efficient days.}$$

- BOTH  $x$  &  $Q^2$ -RANGES ARE EXTENDED BY TWO ORDERS OF MAGNITUDE
  - $x \sim 10^{-4}$
  - $Q^2 \sim 0 (10^4 \text{ GeV}^2)$
- WIDE RANGE IN  $Q^2$  FOR QCD TEST
- HUGE STATISTICS AT LOW  $x$  (NC)
- $e^{\pm} p \rightarrow \bar{\nu}_e X$  MEASURABLE AT HIGH  $Q^2$

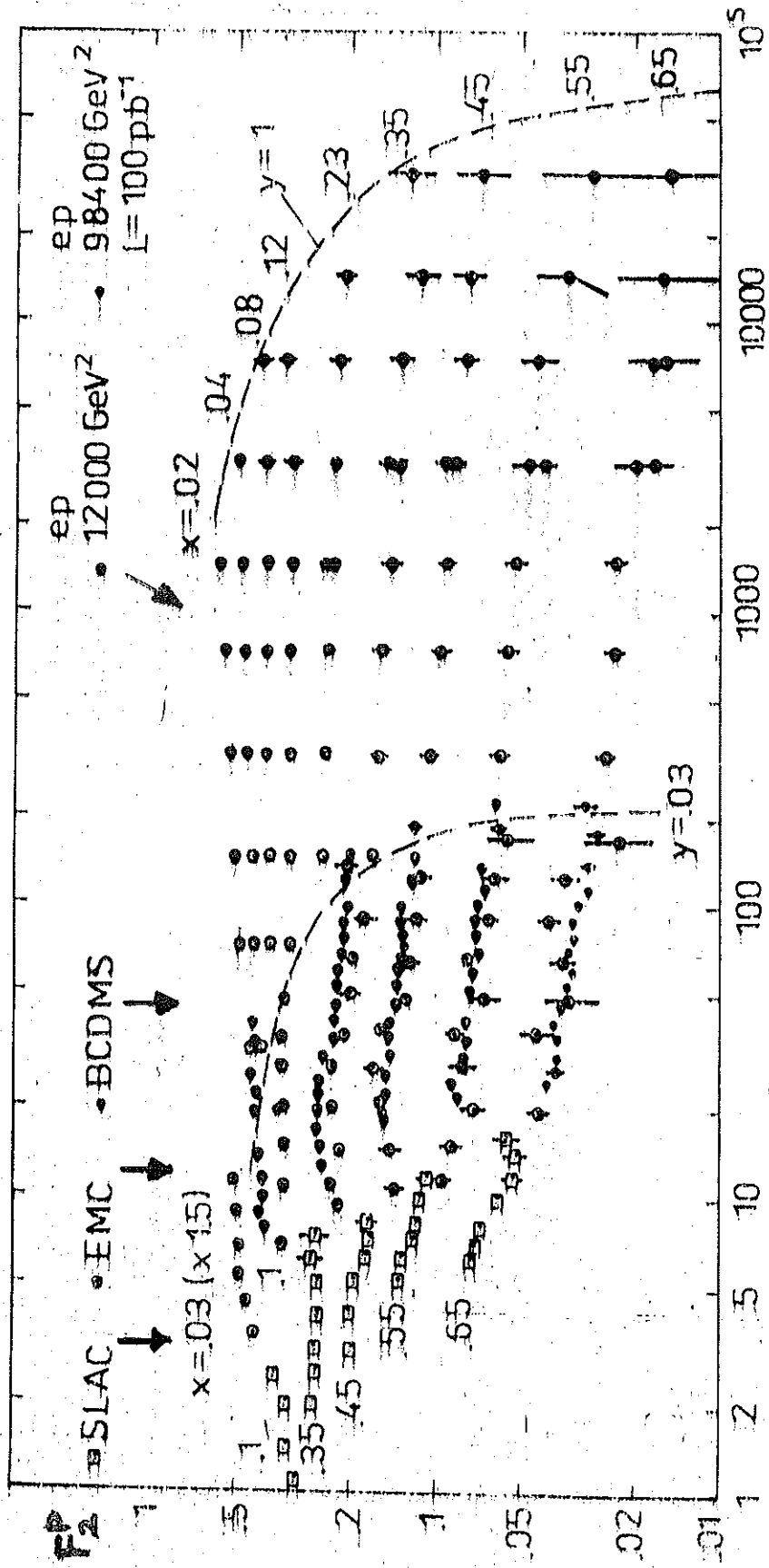
Fig. 1

$\sigma_s$  [GeV<sup>2</sup>s]





# 5 ORDERS OF MAGNITUDE IN $Q^2$



## 2. QUARK DISTRIBUTIONS AT HIGH $Q^2$

FROM 4 CROSS SECTIONS (AT  $A_e=10$ )

$$\frac{d^2\sigma_{NC}^{e^+p}}{dx dy} \quad \frac{d^2\sigma_{NC}^{e^-p}}{dx dy} \quad \frac{d^2\sigma_{CC}^{e^+p}}{dx dy} \quad \frac{d^2\sigma_{CC}^{e^-p}}{dx dy}$$

4 COMBINATIONS OF QUARK-DISTRIBUTIONS CAN BE UNFOLDED IN PRINCIPLE.

- NOTE, HOWEVER, THAT  $\sigma_{NC}^{e^+p}$  &  $\sigma_{NC}^{e^-p}$  ARE ALMOST IDENTICAL AT LOW  $Q^2$  ( $Q^2 \lesssim 1000 \text{ GeV}^2$ )
- THAT A REASONABLE CC-STATISTICS CAN ONLY BE MEASURED AT HIGH  $Q^2$  ( $Q^2 \gtrsim 500 \text{ GeV}^2$ )

TWO METHODS :

(J.B., M. Klein, T. Naumann, T. Riemann, HERA proc. I, p.69 ;  
G. Ingelman, R. Ruckl, Z.Phys. C in press) DESY 89-25

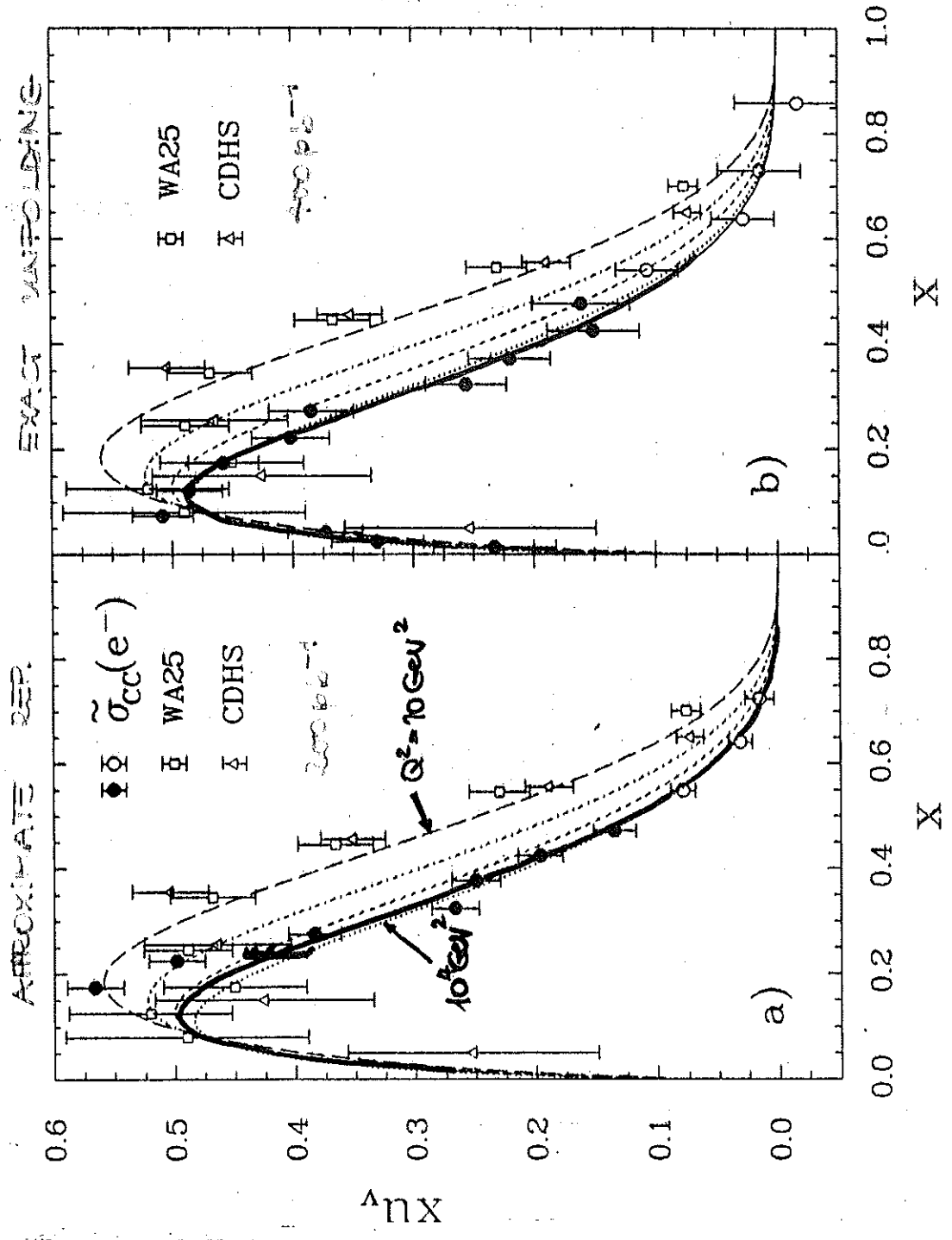
- EXACT UNFOLDING: e.g.

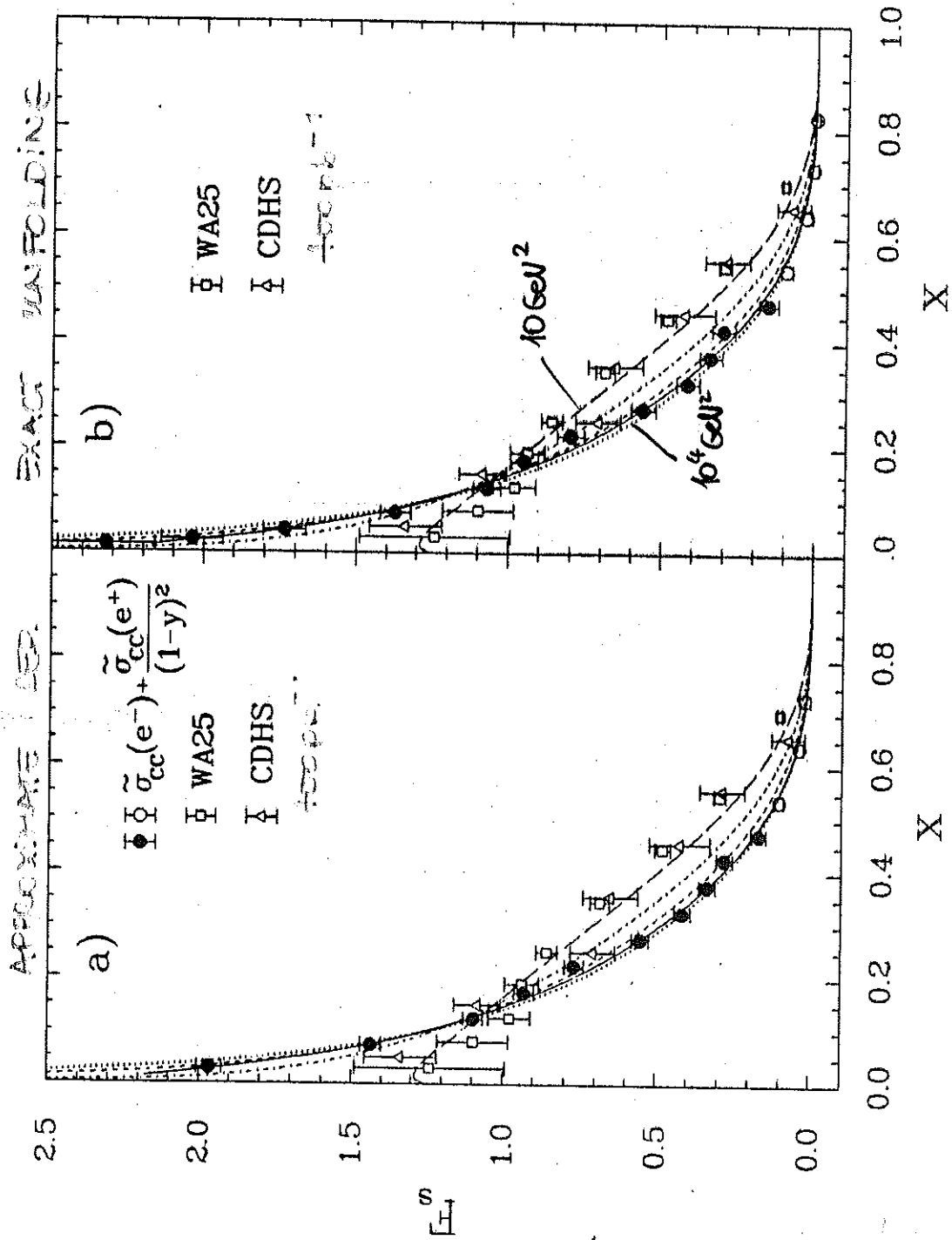
$$\begin{pmatrix} \sigma_{NC}^{e^+p} \\ \sigma_{NC}^{e^-p} \\ \sigma_{CC}^{e^+p} \\ \sigma_{CC}^{e^-p} \end{pmatrix} = \left( A_{ij}(y, Q^2) \right) \begin{pmatrix} u_v \\ d_v \\ \sum_i u_i + \bar{u}_i \\ \sum_i d_i + \bar{d}_i \end{pmatrix}$$

- APPROXIMATE REPRESENTATIONS :

e.g.,  $\sigma_{CC}^{e^-p} \propto x u_v$  ,  $\sigma_{CC}^{e^+p} \propto x d_v$  ,  $x \gtrsim 0.25$

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$$F_2^S = x \sum_i (q_i + \bar{q}_i)$$



### 3. $\Lambda_{QCD}$ & $\alpha_s(Q^2)$

\*)

$$F(x, Q^2) = a F_{NS}(x, Q^2) + b F_S(x, Q^2)$$

$$\frac{\partial}{\partial t} F_{NS}(x, t) = P_{qq}(x) \otimes F_{NS}(x, t) \quad (N)$$

$$\frac{\partial}{\partial t} \begin{pmatrix} F_S(x, t) \\ G(x, t) \end{pmatrix} = \begin{pmatrix} P_{qq}(x) & 2N_f P_{qg}(x) \\ P_{gq}(x) & P_{gg}(x) \end{pmatrix} \otimes \begin{pmatrix} F_S(x, t) \\ G(x, t) \end{pmatrix} \quad (S)$$

$$t = \frac{2}{\beta_0} \ln \left[ \ln(Q^2/\Lambda^2) / \ln(Q_0^2/\Lambda^2) \right] ; \beta_0 = 33 - 2N_f$$

AMONG THE COMBINATIONS OF PARTON DISTRIBUTIONS WHICH CAN BE EXTRACTED AT HERA

$$u_v, d_v, u_v \pm d_v, e_u u_v - e_d d_v, \sum_i (q_i + \bar{q}_i), \sum_i (u_i + \bar{u}_i), \sum_i (d_i + \bar{d}_i), F_2 = \sum_i e_i^2 (q_i + \bar{q}_i) \text{ etc.}$$

ONLY  $F_2$  ALLOWS A MEANINGFUL QCD TEST.

$$F_2(x, Q^2) = \frac{1}{6} x \Delta_p(x, Q^2) + \frac{5}{18} x F_S(x, Q^2)$$

$$\Delta_p = \sum_i (u_i + \bar{u}_i - d_i - \bar{d}_i)$$

$$F_S = \sum_i (q_i + \bar{q}_i)$$

\*) J.B., G. INGELMAN, M. KLEIN, R. RUCKL, PHE 89-01, DESY 89-101.

i) STATISTICAL PRECISION :

MC-DATA :  $\Lambda = 200 \text{ MeV}$

$\sqrt{s}$ [GeV]	$\int \mathcal{L} dt$ [pb <sup>-1</sup> ]	$Q_{min}^2$ [GeV <sup>4</sup> ]	$x_{min}$	FIT	$\Lambda / \text{MeV}$	RESTRICTED RANGE $\Lambda / \text{MeV}$
314	200	100	.25	(N)	$145 \pm 48$	$175 \pm 176$
			.01	(N) + (S)	$297 \pm 76$	$177 \pm 135$
			.01	-1-, $xG(Q_0^2)$ fix	$215 \pm 16$	$201 \pm 25$
	100	10	$10^{-4}$	(N) + (S)	$196 \pm 5$	$225 \pm 25$
134	100	20	.25	(N)	$200 \pm 47$	$460 \pm 263$
			.01	(N) + (S)	$211 \pm 27$	$227 \pm 58$

ii) SYSTEMATIC EFFECTS :

- RESTRICTED RANGE : SKEWING EFFECTS  $\leq 10\%$   
J. FELTESSE
- CALIBRATION UNCERTAINTIES OF THE ELECTROMAGNETIC & HADRONIC CALORIMETER

$$\hat{E}_e = E_e (1 + \epsilon_e) \quad ; \quad \hat{E}_H = E_H (1 + \epsilon_H)$$

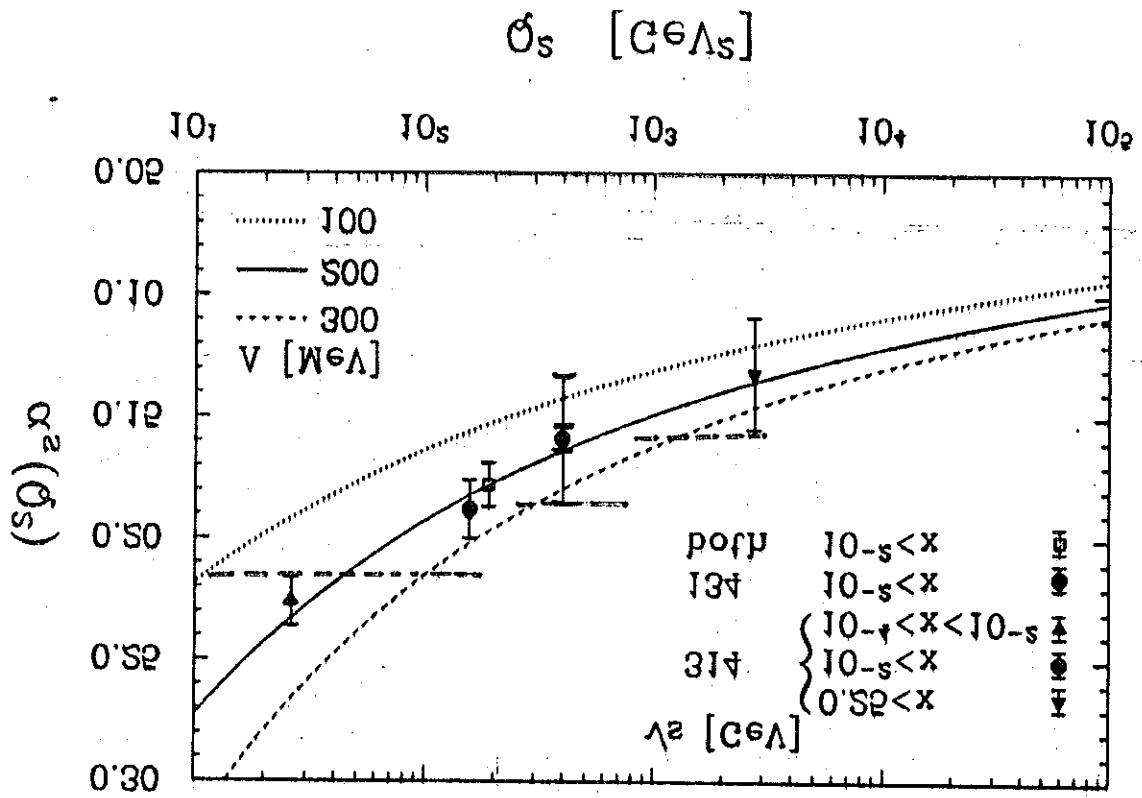
WE CONSIDER : a)  $x \geq 10^{-2}$ ,  $Q^2 \geq 100 \text{ GeV}^2$  : JET-MEASUREMENT  
 b)  $10^{-4} \leq x \leq 10^{-2}$ ,  $Q^2 \leq 100 \text{ GeV}^2$  : ELECTRON-MEASUREMENT

a)  $\Delta\Lambda = \pm 70 \text{ MeV}$ ,  $\epsilon_e = \pm 0.01$   
 b)  $\Delta\Lambda = \pm 45 \text{ MeV}$ ,  $\epsilon_e = \pm 0.01$  } SCALES LINEARLY WITH  $\epsilon$  FOR  $\epsilon \ll 1$ .

STATISTICAL ERRORS AT  $\int \mathcal{L} dt = 200 \text{ pb}^{-1}$  :

- a)  $\delta\Lambda = 135 \text{ MeV}$
- b)  $\delta\Lambda = 25 \text{ MeV}$  (LO QCD)

Fig. 8



# 4. THE GLUON DISTRIBUTION

(d) QCD ANALYSIS :  $xG(x, Q_0^2)$ ,  $Q_0^2 = 4 \text{ GeV}^2$

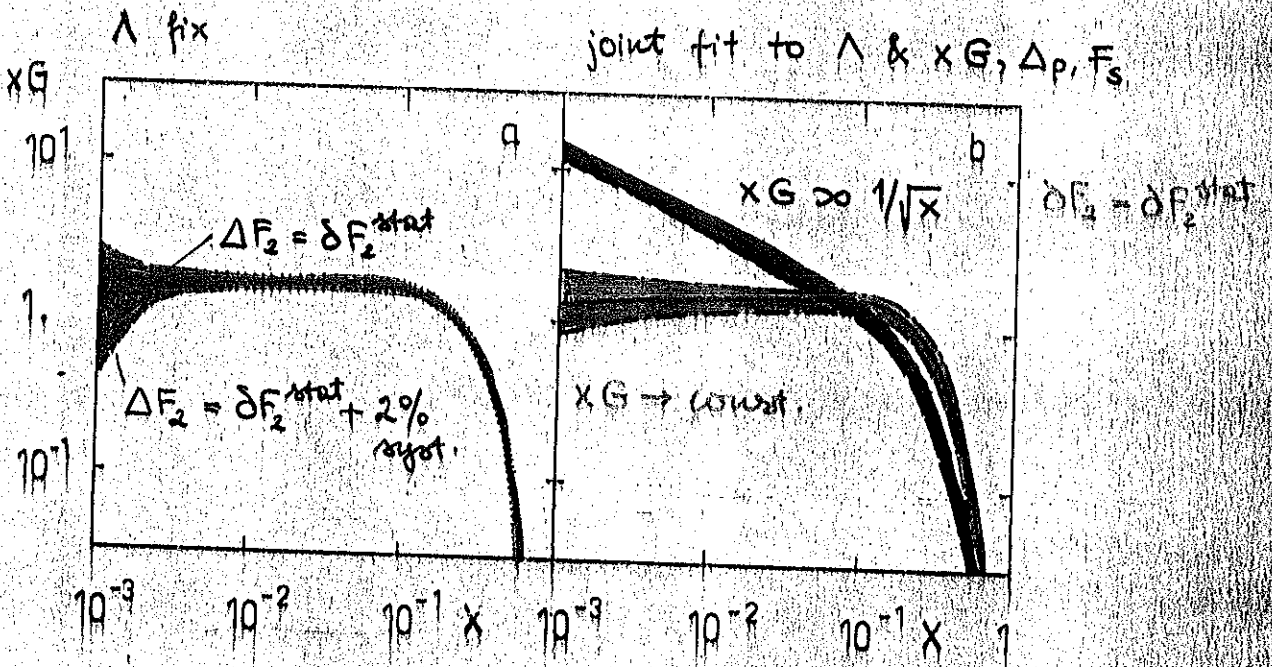
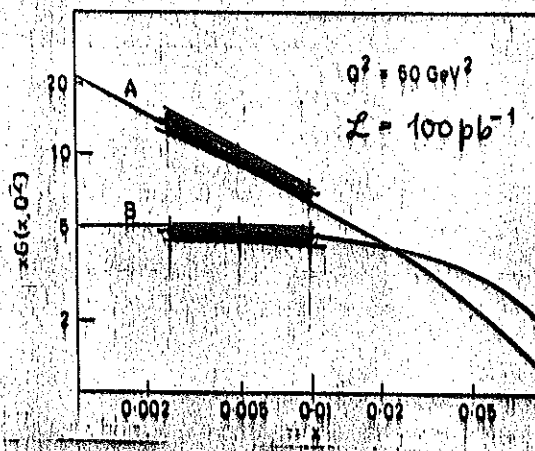
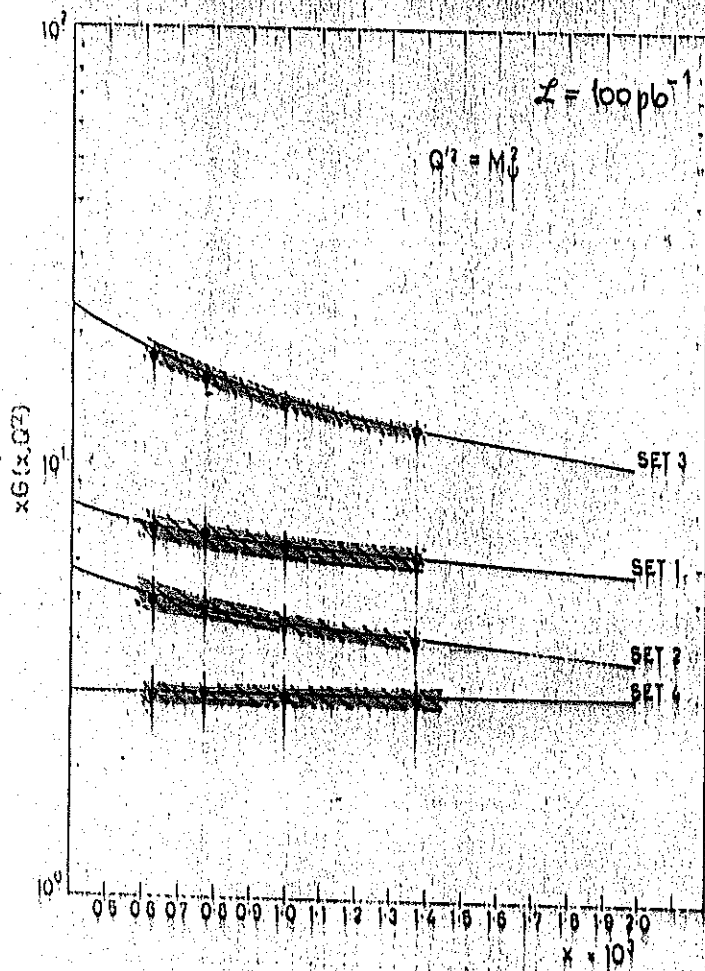


Fig. 4 a)  $xG(\text{stat})$  and  $(\text{stat.} \oplus \text{syst.})$  for  $\delta\Lambda = 0$   
 b)  $xG(\text{stat})$  from joint fit of  $\Lambda$  and  $xG$  with  $xG(\text{DO})$  and  $xG \sim 1/\sqrt{x}$ ;  $x \geq 10^{-3}$ ,  $L = 100 \text{ pb}^{-1}$



A.M. COOPER - SARKAR, et al.  
HERA '88 Vol. 1, p. 231

S.M. THACZYK et al. HERA '88,  
Vol. 1 p. 265

1)  $J/\psi$  - PRODUCTION CROSS SECTION

$$\frac{d\sigma^{ep}_{J/\psi}}{dy} = \frac{15\alpha}{\pi} \frac{1}{y} Y_+ \bar{x} G(\bar{x}, Q^2 = M_{J/\psi}^2) \times \log \frac{Q_{max}^2}{Q_{min}^2} nb$$

$$\bar{x} \sim 3.4 M_{J/\psi}^2 / s_1$$

II) NORMALIZATION UNCERTAINTY (ZEUS)

3)  $\gamma\gamma$  - FUSION (d'Agnostini et al, Services Physics at Future Colliders)

c)  $F_L$  - MEASUREMENT

$$F_L^{QCD}(x, Q^2) = \frac{\alpha_s}{4\pi} \cdot (I_1 + I_2)$$

$$I_1 = \frac{16}{3} \int_x^1 \frac{dy}{y} \left(\frac{x}{y}\right)^2 F_2(y, Q^2)$$

$$I_2 = 8 \sum_i e_i^2 \int_x^1 \frac{dy}{y} \left(\frac{x}{y}\right)^2 (1 - \frac{x}{y}) \cdot G(y, Q^2)$$

## 5. CONCLUSIONS

1. HERA WILL PROBE THE PROTON STRUCTURE VIA DEEP INELASTIC SCATTERING
  - UP TO  $Q^2 \sim O(10^4 \text{ GeV}^2)$
  - DOWN TO  $x \sim 10^{-4}$
2. QUARK DISTRIBUTIONS CAN BE UNFOLDED FROM CROSS SECTIONS AS FUNCTIONS OF  $x$  FOR  $Q^2 \sim 10^3 \dots 10^4 \text{ GeV}^2$ .
3. THE LOW  $x$  RANGE  $10^{-4} < x < 10^{-2}$  FOR  $10 < Q^2 < 10^3 \text{ GeV}^2$  COULD BECOME A NEW TESTING GROUND FOR QCD.  $F_2(Q^2, x)$  CAN BE MEASURED THERE WITH HIGH STATISTICS. THERE ARE STILL MORE EFFORTS NEEDED TO WORK OUT THE THEORY OF QCD EVOLUTION IN THIS RANGE.
4. THE GLUON DISTRIBUTION CAN BE DETERMINED BY DIFFERENT METHODS, PARTICULARLY AT LOW  $x$ .
5. THE PRECISE MEASUREMENTS OF THE ABOVE QUANTITIES ARE LONG TERM TASKS REQUIRING IN THE CASE OF THE QCD TEST TO CONTROL SYSTEMATICS AT THE PER CENT LEVEL.