

Leptoquark Pair Production

in $\gamma\gamma$ Scattering

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DESY

1. Introduction
2. Classification of Leptoquark States
3. $\gamma\gamma \rightarrow LQ\bar{L}\bar{Q}$
4. QCD Corrections
5. Conclusions

1. Classification of Leptoquark States

B and L conserving

family-diagonal

BUCHMÜLLER,
RÖCKL, WYLER
'87

$SU(3)_c \times SU(2)_L \times U(1)_Y$ invariant couplings

leptoquark (Φ)	spin	F	colour	T_3	Q_{em}	$\lambda_L(lq)$	$\lambda_R(lq)$	$\lambda_L(\nu q)$
S_1	0	-2	$\bar{3}$	0	1/3	g_{1L}	g_{1R}	$-g_{1L}$
\tilde{S}_1	0	-2	$\bar{3}$	0	4/3	0	\tilde{g}_{1R}	0
\tilde{S}_3	0	-2	$\bar{3}$	+1	4/3	$-\sqrt{2}g_{3L}$	0	0
				0	1/3	$-g_{3L}$	0	$-g_{3L}$
R_2	0	0	3	1/2	5/3	h_{2L}	h_{2R}	0
				-1/2	2/3	0	$-h_{2R}$	h_{2L}
\tilde{R}_2	0	0	3	1/2	2/3	\tilde{h}_{2L}	0	0
				-1/2	-1/3	0	0	\tilde{h}_{2L}
$V_{2\mu}$	1	-2	$\bar{3}$	1/2	4/3	g_{2L}	g_{2R}	0
				-1/2	1/3	0	g_{2R}	g_{2L}
$\tilde{V}_{2\mu}$	1	-2	$\bar{3}$	1/2	1/3	\tilde{g}_{2L}	0	0
				-1/2	-2/3	0	0	\tilde{g}_{2L}
$U_{1\mu}$	1	0	3	0	2/3	h_{1L}	h_{1R}	h_{1L}
$\tilde{U}_{1\mu}$	1	0	3	0	5/3	0	h_{1R}	0
$\tilde{U}_{3\mu}$	1	0	3	+1	5/3	$\sqrt{2}h_{3L}$	0	0
				0	2/3	$-h_{3L}$	0	h_{3L}
				-1	-1/3	0	0	$\sqrt{2}h_{3L}$

$$\mathcal{L} = \mathcal{L}_{|F|=2}^f + \mathcal{L}_{F=0}^f + \mathcal{L}^{\gamma, Z, g}$$

$$\mathcal{L}^{\gamma, Z, g} = \sum_{\text{scalars}} [(D^\mu \Phi)^\dagger (D_\mu \Phi) - M^2 \Phi^\dagger \Phi] + \sum_{\text{vectors}} \left[-\frac{1}{2} G_{\mu\nu}^\dagger G^{\mu\nu} + M^2 \Phi^\dagger \Phi_\mu \right]$$

$$D_\mu = \partial_\mu - ieQ^A A_\mu - ieQ^Z Z_\mu - ig_s \frac{\lambda_a}{2} A_\mu^a$$

$$\begin{aligned} \mathcal{L}_{F=0}^f = & (h_{2L} \bar{u}_R l_L + h_{2R} \bar{q}_L i\tau_2 e_R) R_2 + \tilde{h}_{2L} \bar{d}_R l_L \tilde{R}_2 \\ & + (h_{1L} \bar{q}_L \gamma^\mu l_L + h_{1R} \bar{d}_R \gamma^\mu e_R) U_{1\mu} \\ & + \tilde{h}_{1R} \bar{u}_R \gamma^\mu e_R \tilde{U}_{1\mu} + h_{3L} \bar{q}_L \vec{\tau} \gamma^\mu l_L \vec{U}_{3\mu} + h.c. \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{|F|=2}^f = & (g_{1L} \bar{q}_L^c i\tau_2 l_L + g_{1R} \bar{u}_R^c e_R) S_1 \\ & + \tilde{g}_{1R} \bar{d}_R^c e_R \tilde{S}_1 + g_{3L} \bar{q}_L^c i\tau_2 \vec{\tau} l_L \vec{S}_3 \\ & + (g_{2L} \bar{d}_R^c \gamma^\mu l_L + g_{2R} \bar{q}_L^c \gamma^\mu e_R) V_{2\mu} \\ & + \tilde{g}_{2L} \bar{u}_R^c \gamma^\mu l_L \tilde{V}_{2\mu} + h.c., \end{aligned}$$

PRESENT LQ BOUNDS

95% CL

1ST GENERATION

$$m > 242 \text{ GeV}$$

2ND GENERATION

$$m > 202 \text{ GeV}$$

3RD GENERATION

$$m > \begin{matrix} 93 \text{ GeV} \\ 99 \text{ GeV} \end{matrix} \quad (|Q| = \frac{1}{3})$$

TEVATRON



SEARCH CONTINUES

- TEVATRON
- LHC

} BEFORE TESLA STARTS

MEASURABLE:

κ_e, λ_e

BUT NOT: $\kappa_\gamma, \lambda_\gamma$

Decay Pattern for Pair Production

states	$l^+l^- + 2jets$	$l\nu + 2jets$	$\nu\bar{\nu} + 2jets$
$S_1 \quad U_1$	$\frac{4}{9} \quad 1 \quad \frac{1}{4}$	$\frac{4}{9} \quad 0 \quad \frac{1}{2}$	$\frac{1}{9} \quad 0 \quad \frac{1}{4}$
$R_2^{2/3} \quad V_2^{1/3}$	$\frac{1}{4} \quad 1 \quad 0$	$\frac{1}{2} \quad 0 \quad 0$	$\frac{1}{4} \quad 0 \quad 1$
$S_3^{1/3} \quad U_3^{2/3}$	$\frac{1}{4}$	$\frac{1}{2}$	$\frac{1}{4}$
$\tilde{S}_1 \quad S_3^{4/3} \quad R_2^{5/3} \quad \tilde{R}_2^{2/3}$ $V_2^{4/3} \quad \tilde{V}_2^{1/3} \quad \tilde{U}_1 \quad U_3^{5/3}$	1	0	0
$S_3^{-2/3} \quad \tilde{R}_2^{-1/3} \quad \tilde{V}_2^{-2/3} \quad U_3^{-1/3}$	0	0	1

Table 3: Branching ratios for final states arising from the decays of leptoquarks associated with the first ($l = e$) and second ($l = \mu$) family. The sequence of branching fractions given in the second and third row refers to the assumptions $\lambda_L = \lambda_R$, $\lambda_L = 0$, and $\lambda_R = 0$, respectively.

LINEAR COLLIDER:

$$\begin{array}{l} e^+e^- \rightarrow \phi\bar{\phi} \\ \gamma\gamma \rightarrow \phi\bar{\phi} \end{array} \left. \vphantom{\begin{array}{l} e^+e^- \\ \gamma\gamma \end{array}} \right\} \text{WIDELY INDEPENDENT OF } \lambda_{eq}$$

→ GAUGE COUPLINGS

e^+e^- ANNIHILATION:

HEWETT, RIZZO, ZERWAS, JB RÜCKL
JB BOOS KRYUKOV

MASS BOUNDS: $m_\phi \lesssim \frac{1}{2}\sqrt{s}$

LOSSES: BEAMSTRAHLUNG, QED JB '93
 $-\Delta\sigma \sim 40\%$ $O(\alpha^2)$
needed

ENHANCEMENT: QCD FS $\frac{1}{\beta}$ JB '93

→ NEARLY BALANCED UP TO $\sqrt{s} \sim 1\text{TeV}$.

2. $\gamma\gamma$ Scattering

$U_{em}(1)$ invariant Lagrangian:

$$\mathcal{L} = \mathcal{L}_s + \mathcal{L}_v \quad (1)$$

with

$$\mathcal{L}_s = \sum_{\text{scalars}} \left[(D^\mu \Phi)^\dagger (D_\mu \Phi) - M_s^2 \Phi^\dagger \Phi \right] \quad (2)$$

and

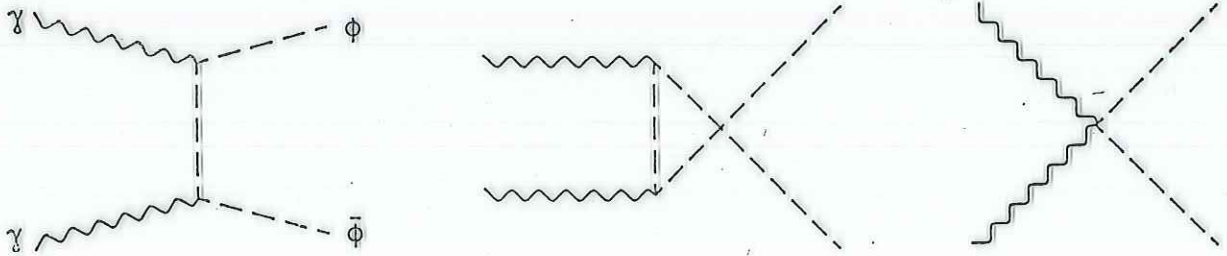
$$\mathcal{L}_v = \sum_{\text{vectors}} \left\{ -\frac{1}{2} G_{\mu\nu}^\dagger G^{\mu\nu} + M_v^2 \Phi_\mu^\dagger \Phi^\mu - ie \left[(1 - \kappa_A) \Phi_\mu^\dagger \Phi_\nu F^{\mu\nu} + \frac{\lambda_A}{M_v^2} G_{\sigma\mu}^\dagger G_\nu^\mu F^{\nu\sigma} \right] \right\} \quad (3)$$

Field strength tensors:

$$\begin{aligned} F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu, \\ G_{\mu\nu} &= D_\mu \Phi_\nu - D_\nu \Phi_\mu \end{aligned} \quad (4)$$

Covariant derivative:

$$D_\mu = \partial_\mu - ieQ_\gamma A_\mu \quad (5)$$



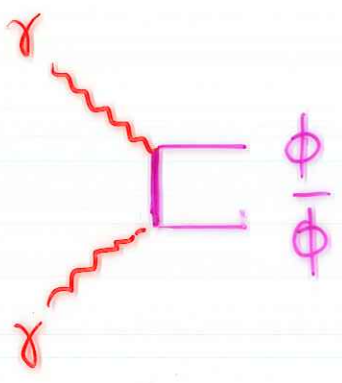
Anomalous couplings: κ_A and λ_A



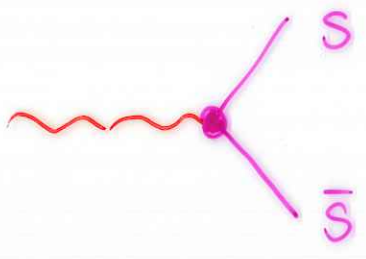
anomalous magnetic moment μ_Φ and electric quadrupole moment q_Φ of leptoquarks:

$$\begin{aligned} \mu_{\Phi,A} &= \frac{eQ_\gamma}{2M_\Phi} (2 - \kappa_A + \lambda_A) \\ q_{\Phi,A} &= -\frac{eQ_\gamma}{M_\Phi^2} (1 - \kappa_A - \lambda_A) \end{aligned} \quad (6)$$

WHY TO SEARCH FOR $\phi\bar{\phi}$ IN $\gamma\gamma$ -SCATTERING?



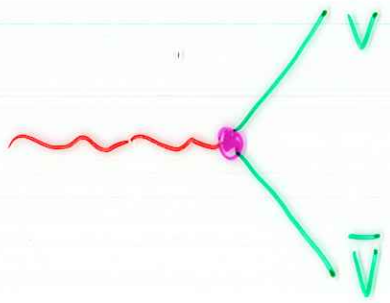
SCALARS & VECTORS



$$\sigma \sim Q_\phi^4$$

$$|Q_\phi| = \frac{1}{3} \dots \frac{5}{3} !$$

$$\sigma = O(1) \dots O(625).$$



$$\sigma_V = \sigma_V(k_\gamma, \lambda_\gamma)$$

ANOMALOUS COUPLINGS.

Sensitivity to Quantum Numbers

Process	LQ	Quantum Numbers
e^+e^- Annihilation	S,V	$Q_\Phi^\gamma, Q_\Phi^Z, \lambda_{L,R}$
$\gamma\gamma$ Collider	S	Q_Φ^γ
	V	$Q_\Phi^\gamma, \kappa_A, \lambda_A$
$e\gamma$ Collider	S,V	$\lambda_{L,R}, Q_\Phi^\gamma$
ep Collider	single LQ	$\lambda_{L,R}$
	pairs : S	Q_Φ^γ
	pairs : V	$Q_\Phi^\gamma, \kappa_A, \kappa_G, \lambda_A, \lambda_G$
$p\bar{p}$ Collider	V	κ_G, λ_G

$\sigma \sim Q_\Phi^4$

2. CROSS SECTIONS

BORN:

J. BLÜMWEIN, E. BOOS, NUCL. PHYS (P.S.) 37B (1994) 181

J. BLÜMWEIN, E. BOOS, A. KRYUKOV, Z. PHYS. C76 (1997) 137

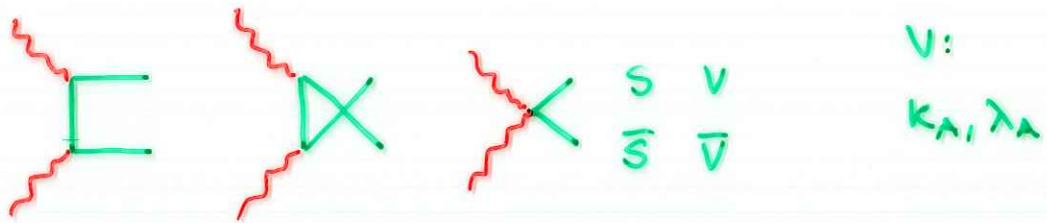
CODE: LQPAIR 1.0 : J. BLÜMWEIN, E. BOOS, A. KRYUKOV
199.

→ CONVOLUT WITH COMPTON LASER SPECTRA

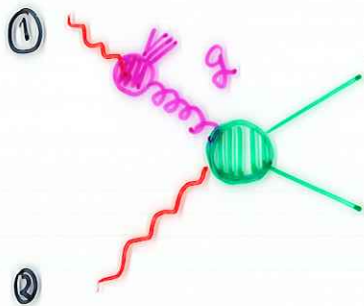
THREE CONTRIBUTIONS:

→	γ_{DIR}	x	γ_{DIR}
MOST IMPORTANT TERM	γ_{DIR}	x	γ_{RESOLVED} ← HADRONIC.
	γ_{RESOLVED}	x	γ_{RESOLVED}

DIRECT TERMS :



DIRECT - RESOLVED CONTRIBUTIONS :

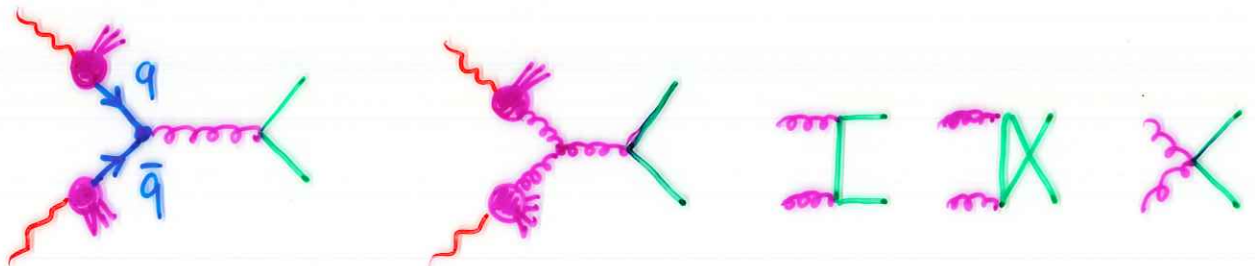


assume: $\lambda_{qe} \ll 1$.

+ ① \leftrightarrow ②.

$V: \underline{\delta}, \underline{g}$
 κ_A, κ_G
 λ_A, λ_G

RESOLVED - RESOLVED TERM :



$V: \kappa_G, \lambda_G$

V: direct.

$$\begin{aligned}
\tilde{F}_0 &= \beta \left(\frac{11}{2} - \frac{9}{4}\beta^2 + \frac{3}{4}\beta^4 \right) - \frac{3}{8} (1 - \beta^2 - \beta^4 + \beta^6) \ln \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_1 &= -8\beta - \frac{3}{2} (1 - \beta^2) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_2 &= 3\beta + \frac{1}{4}\beta \frac{\hat{s}}{M_\Phi^2} + \left(\frac{7}{2} - 2\beta^2 \right) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_3 &= -\frac{1}{4}\beta \frac{\hat{s}}{M_\Phi^2} + \left(-2 + \frac{3}{4}\beta^2 \right) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_4 &= -\frac{1}{96}\beta + \frac{5}{48}\beta \frac{\hat{s}}{M_\Phi^2} + \frac{4 - \beta^2}{16} \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_5 &= - (1 - \beta^2) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_6 &= -\frac{1}{6}\beta + \frac{17}{12}\beta \frac{\hat{s}}{M_\Phi^2} + \left(-3 - \frac{\beta^2}{2} + \frac{1}{2} \frac{\hat{s}}{M_\Phi^2} \right) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_7 &= -\beta + \frac{11}{6}\beta \frac{\hat{s}}{M_\Phi^2} - \frac{1}{3}\beta \frac{\hat{s}^2}{M_\Phi^4} - \frac{3 + \beta^2}{4} \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_8 &= -\frac{1}{96}\beta + \frac{59}{80}\beta \frac{\hat{s}}{M_\Phi^2} - \frac{113}{320}\beta \frac{\hat{s}^2}{M_\Phi^4} + \frac{43}{960}\beta \frac{\hat{s}^3}{M_\Phi^6} + \left(-\frac{1}{2} - \frac{1}{16}\beta^2 + \frac{1}{8} \frac{\hat{s}}{M_\Phi^2} \right) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_9 &= 2\beta + (2 + \beta^2) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_{10} &= 2\beta - \frac{7}{3}\beta \frac{\hat{s}}{M_\Phi^2} + \left(3 + \frac{5}{4}\beta^2 - \frac{1}{2} \frac{\hat{s}}{M_\Phi^2} \right) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_{11} &= \frac{1}{24}\beta - \frac{59}{48}\beta \frac{\hat{s}}{M_\Phi^2} + \frac{5}{32}\beta \frac{\hat{s}^2}{M_\Phi^4} + \frac{5 + \beta^2}{4} \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_{12} &= -\beta - \frac{1}{2}\beta \frac{\hat{s}}{M_\Phi^2} + \left(-\frac{1}{4} - \frac{7}{4}\beta^2 \right) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_{13} &= \frac{1}{24}\beta + \frac{1}{3}\beta \frac{\hat{s}}{M_\Phi^2} - \frac{1}{4} (1 - \beta^2) \log \left| \frac{1+\beta}{1-\beta} \right| \\
\tilde{F}_{14} &= -\frac{1}{16}\beta + \frac{11}{96}\beta \frac{\hat{s}}{M_\Phi^2} + \frac{17}{192}\beta \frac{\hat{s}^2}{M_\Phi^4} + \left(\frac{1}{8} \frac{\hat{s}}{M_\Phi^2} - \frac{3}{4} - \frac{3}{8}\beta^2 \right) \log \left| \frac{1+\beta}{1-\beta} \right| \tag{12}
\end{aligned}$$

Tree-level unitarity:

$$\lambda_A = 0$$

$$\kappa_A^2 \left[\left(\kappa_A - \frac{6}{5} \right)^2 + \frac{24}{25} \right] = 0. \tag{13}$$

Since κ_A, λ_A real $\rightarrow \kappa_A \equiv \lambda_A \equiv 0$.

$$F_{18} = 2(5 - \beta^2 \cos^2 \theta) - \frac{\hat{s}}{M_\Phi^2} \frac{11 - 15\beta^2 \cos^2 \theta + 4\beta^4 \cos^4 \theta}{4} - \frac{\hat{s}^2 (1 - \beta^2 \cos^2 \theta)^2}{4M_\Phi^4},$$

$$F_{19} = 3 - \beta^2 \cos^2 \theta - \frac{\hat{s}}{M_\Phi^2} \frac{7 - 8\beta^2 \cos^2 \theta + \beta^4 \cos^4 \theta}{4} + \frac{\hat{s}^2}{M_\Phi^4} \frac{11 - 13\beta^2 \cos^2 \theta + \beta^4 \cos^4 \theta + \beta^6 \cos^6 \theta}{32} + \frac{\hat{s}^3}{M_\Phi^6} \frac{5 - 7\beta^2 \cos^2 \theta - \beta^4 \cos^4 \theta + 3\beta^6 \cos^6 \theta}{128},$$

$$F_{20} = -\frac{3 - \beta^2 \cos^2 \theta}{2} + \frac{\hat{s}}{M_\Phi^2} \frac{(1 - \beta^2 \cos^2 \theta)^2}{8} + \frac{\hat{s}^2}{M_\Phi^4} \frac{11 - 23\beta^2 \cos^2 \theta + 13\beta^4 \cos^4 \theta - \beta^6 \cos^6 \theta}{64}.$$

The functions $\bar{F}_i(\hat{s}, \beta)$, which describe the different contributions to the integrated cross-section (13), are:

$$\bar{F}_0 = \beta \left(\frac{11}{2} - \frac{9}{4}\beta^2 + \frac{3}{4}\beta^4 \right) - \frac{3}{8} (1 - \beta^2 - \beta^6) \ln \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_1 = -4\beta - \frac{3}{4}(1 - \beta^2) \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_2 = \frac{1}{16} \beta \frac{\hat{s}}{M_\Phi^2} + \frac{3 - \beta^2}{4} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_3 = 3\beta + \frac{1}{8} \beta \frac{\hat{s}}{M_\Phi^2} + \left(2 - \frac{3}{2}\beta^2 \right) \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_4 = -\frac{1}{8} \beta \frac{\hat{s}}{M_\Phi^2} + \left(-1 + \frac{3}{8}\beta^2 \right) \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_5 = -\frac{1}{96} \beta + \frac{5}{48} \beta \frac{\hat{s}}{M_\Phi^2} + \frac{4 - \beta^2}{16} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_6 = -\frac{1}{2}(1 - \beta^2) \log \left| \frac{1 + \beta}{1 - \beta} \right|.$$

$$\bar{F}_7 = \frac{7}{12} \beta \frac{\hat{s}}{M_\Phi^2} + \frac{1}{24} \beta \frac{\hat{s}^2}{M_\Phi^4} - \frac{5 + \beta^2}{4} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_8 = -\frac{1}{6} \beta + \frac{1}{4} \beta \frac{\hat{s}}{M_\Phi^2} - \frac{1}{12} \beta \frac{\hat{s}^2}{M_\Phi^4} + \left(-\frac{1}{2} + \frac{1}{2} \beta \frac{\hat{s}}{M_\Phi^2} \right) \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_9 = -\frac{1}{2} \beta + \frac{11}{12} \beta \frac{\hat{s}}{M_\Phi^2} - \frac{1}{6} \beta \frac{\hat{s}^2}{M_\Phi^4} - \frac{3 + \beta^2}{8} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{10} = -\frac{1}{96} \beta + \frac{59}{80} \beta \frac{\hat{s}}{M_\Phi^2} - \frac{113}{320} \beta \frac{\hat{s}^2}{M_\Phi^4} + \frac{43}{960} \beta \frac{\hat{s}^3}{M_\Phi^6} + \left(-\frac{1}{2} - \frac{1}{16} \beta^2 + \frac{1}{8} \beta \frac{\hat{s}}{M_\Phi^2} \right) \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{11} = \frac{1}{2}(1 + \beta^2) \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{12} = \beta + \frac{1}{2} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{13} = \beta - \frac{5}{12} \beta \frac{\hat{s}}{M_\Phi^2} + \frac{1}{24} \beta \frac{\hat{s}^2}{M_\Phi^4} + \left[-\frac{1}{4} \beta \frac{\hat{s}}{M_\Phi^2} + \left(\frac{3}{8} + \frac{1}{4} \beta^2 \right) \log \left| \frac{1 + \beta}{1 - \beta} \right| \right], \quad (\text{A.2})$$

$$\bar{F}_{14} = -\frac{11}{24} \beta \frac{\hat{s}}{M_\Phi^2} - \frac{1}{24} \beta \frac{\hat{s}^2}{M_\Phi^4} + \frac{9 + 3\beta^2}{8} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{15} = \frac{1}{48} \beta - \frac{59}{96} \beta \frac{\hat{s}}{M_\Phi^2} + \frac{5}{64} \beta \frac{\hat{s}^2}{M_\Phi^4} + \frac{5 + \beta^2}{8} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{16} = -\frac{1}{2} \beta - \frac{1}{8} \beta^2 \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{17} = -\frac{1}{96} \beta + \frac{1}{48} \beta \frac{\hat{s}}{M_\Phi^2} + \frac{1}{48} \beta \frac{\hat{s}^2}{M_\Phi^4} - \frac{2 + \beta^2}{16} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

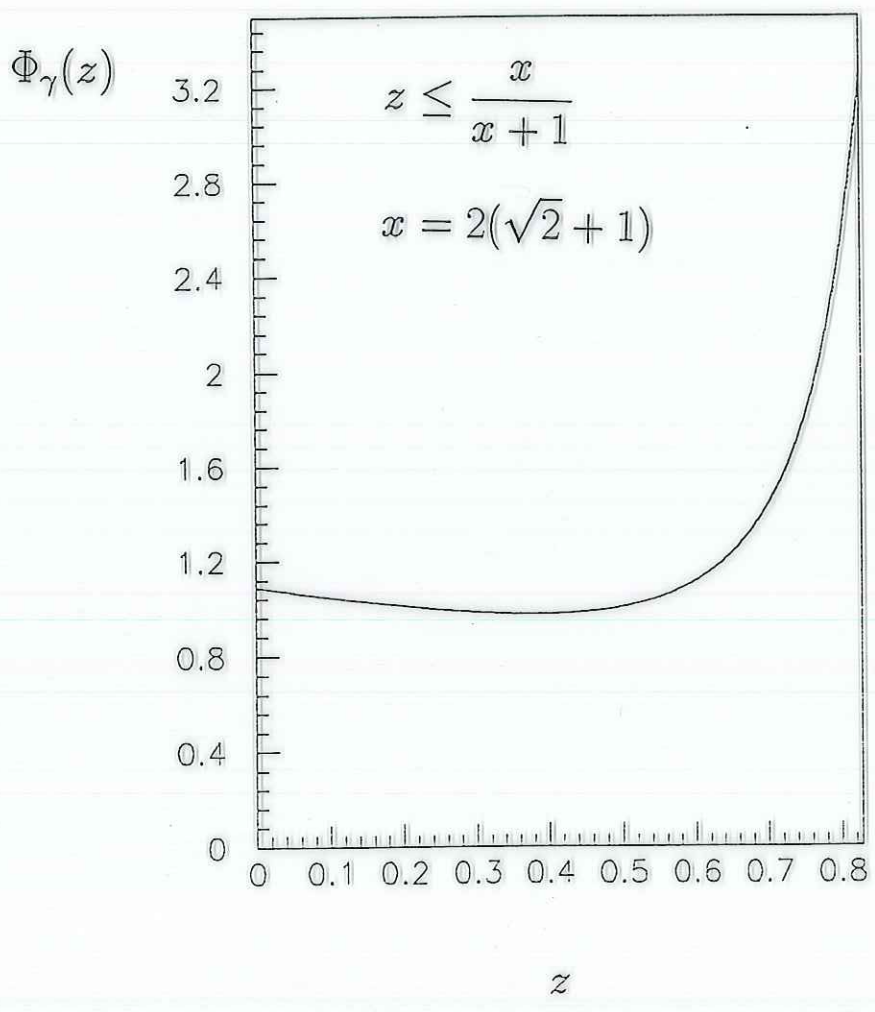
$$\bar{F}_{18} = -\frac{1}{4} \beta \frac{\hat{s}}{M_\Phi^2} - \frac{1 - 6\beta^2}{8} \log \left| \frac{1 + \beta}{1 - \beta} \right|,$$

$$\bar{F}_{19} = -\frac{1}{24} \beta + \frac{7}{96} \beta \frac{\hat{s}}{M_\Phi^2} + \frac{3}{64} \beta \frac{\hat{s}^2}{M_\Phi^4} + \left[\frac{1}{8} \beta \frac{\hat{s}}{M_\Phi^2} - \frac{2 + \beta^2}{4} \log \left| \frac{1 + \beta}{1 - \beta} \right| \right],$$

$$\bar{F}_{20} = \frac{1}{48} \beta + \frac{1}{6} \beta \frac{\hat{s}}{M_\Phi^2} - \frac{1}{8} (1 - \beta^2) \log \left| \frac{1 + \beta}{1 - \beta} \right|.$$

$$\Phi_\gamma(z) = \frac{1}{N(x)} \left[1 - z + \frac{1}{1-z} - \frac{4z}{x(1-z)} + \frac{4z^2}{x^2(1-z)^2} \right]$$

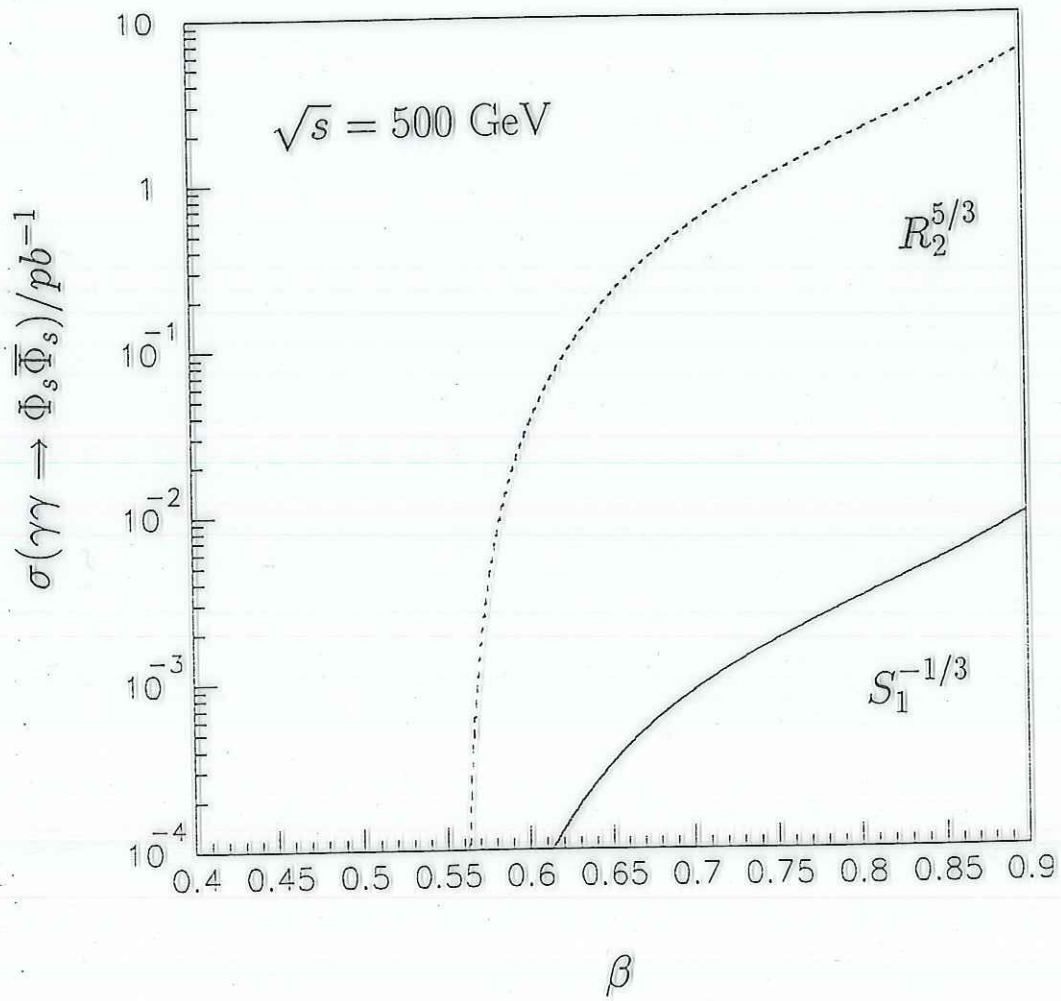
$$N(x) = \frac{16 + 32x + 18x^2 + x^3}{2x(1+x)^2} + \frac{x^2 - 4x - 8}{x^2} \ln(1+x)$$

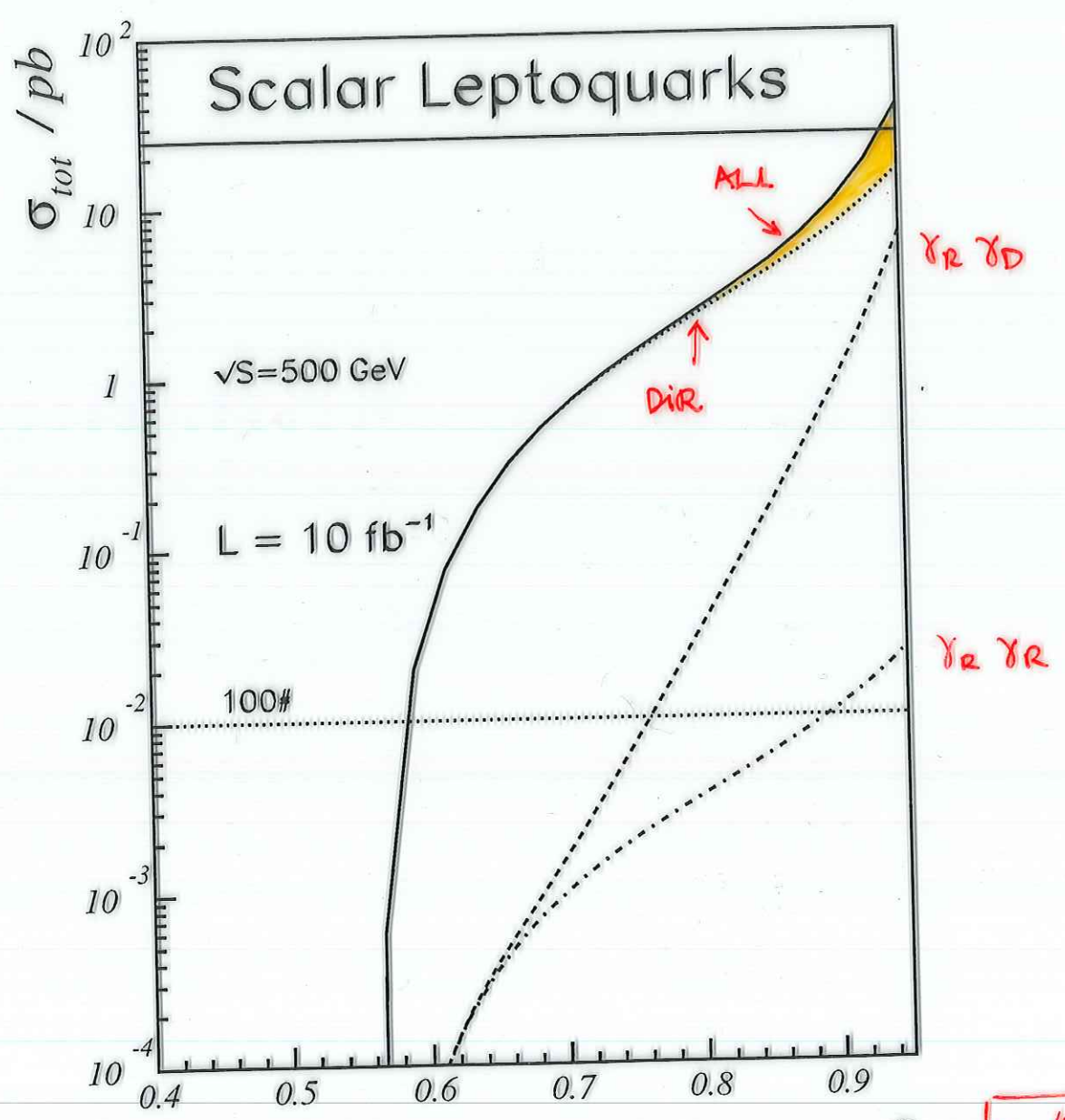


COMPTON -
CONVERSION
SPECTRUM

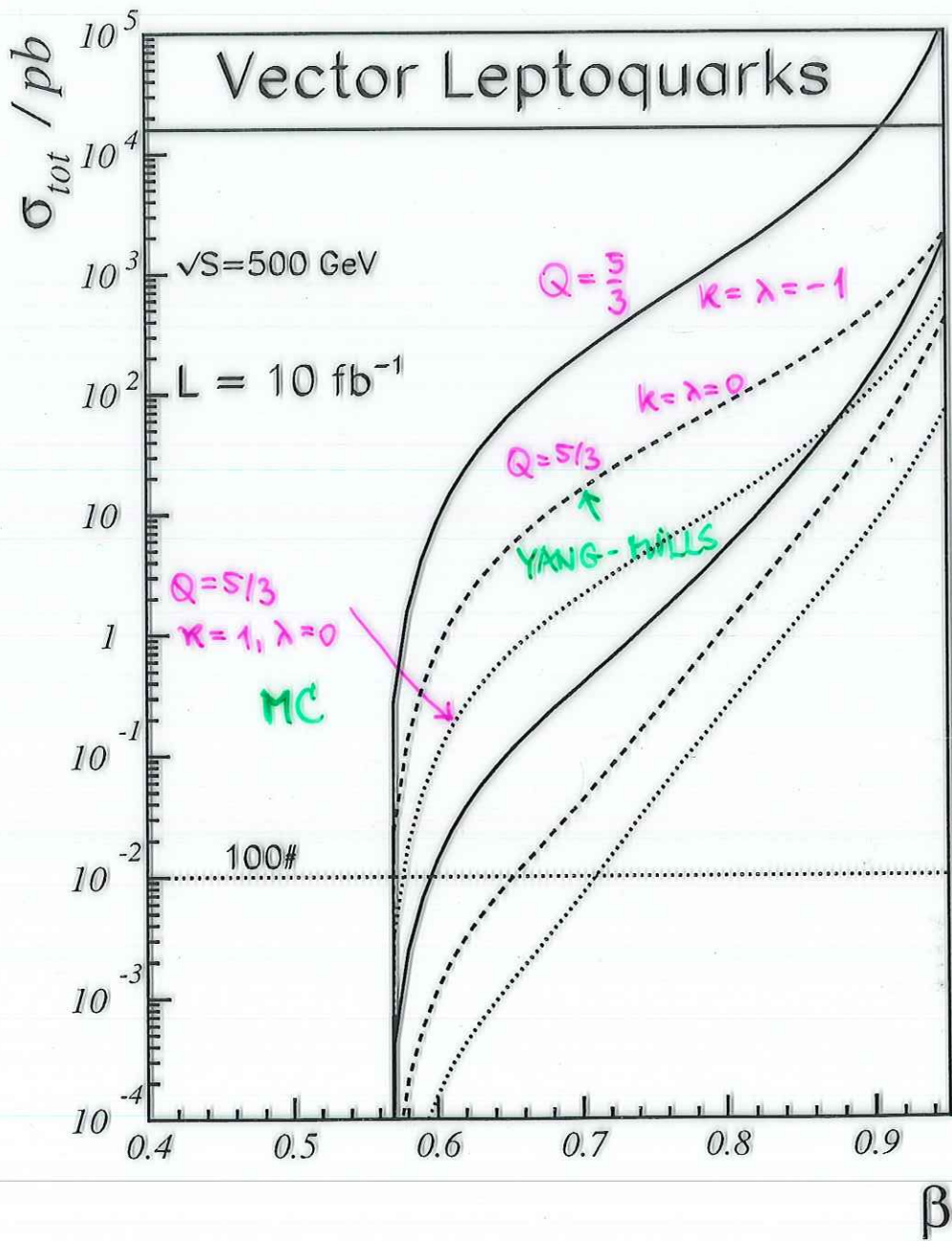
$$\sigma_{scalar}(s) = \frac{\pi\alpha^2}{s} Q_\Phi^4 \left\{ 2(2 - \beta^2)\beta - (1 - \beta^4) \ln \left| \frac{1 + \beta}{1 - \beta} \right| \right\}$$

$$\sigma = \int_0^{z_{max}} dz_1 \int_0^{z_{max}} dz_2 \Phi_\gamma(z_1) \Phi_\gamma(z_2) \hat{\sigma}(\hat{s}) \theta(\hat{s} - 4M_\Phi^2)$$





$$\beta = \sqrt{1 - \frac{4M^2}{s}}$$



BIGI, GABBIANI '91

FADIN, KHOZE '87, '88
BIGI, FADIN, KHOZE '91
PESKIN, STRASSLER '91.

$$d\sigma = d\sigma_B \cdot \frac{K_S(E)}{K_S^{(B)}(E)}$$

$$K_{(B)}(E) = \frac{M_S}{4\pi} \sqrt{M_S E}$$

$$K_S(E) = \frac{M_S^2}{4\pi} \left\{ \frac{k_+}{M_S} + \frac{2k_1}{M_S} \arctan\left(\frac{k_+}{k_-}\right) \right.$$

$$\left. + \sum_{n=1}^{\infty} \frac{2k_1^2}{M_S^2 n^4} \frac{\Gamma_S k_1 \cdot n + K_+ [n^2 \sqrt{E^2 + \Gamma_S^2} + \frac{k_1^2}{M_S}]}{[E + (k_1^2/M_S^2 n^2)]^2 + \Gamma_S^2} \right\}$$

$$k_1 = \frac{2}{3} \alpha_s \cdot M_S$$

$$k_{\pm} = \sqrt{\frac{M_S}{2} [\sqrt{E^2 + \Gamma_S^2} \pm E]}$$

We assumed:

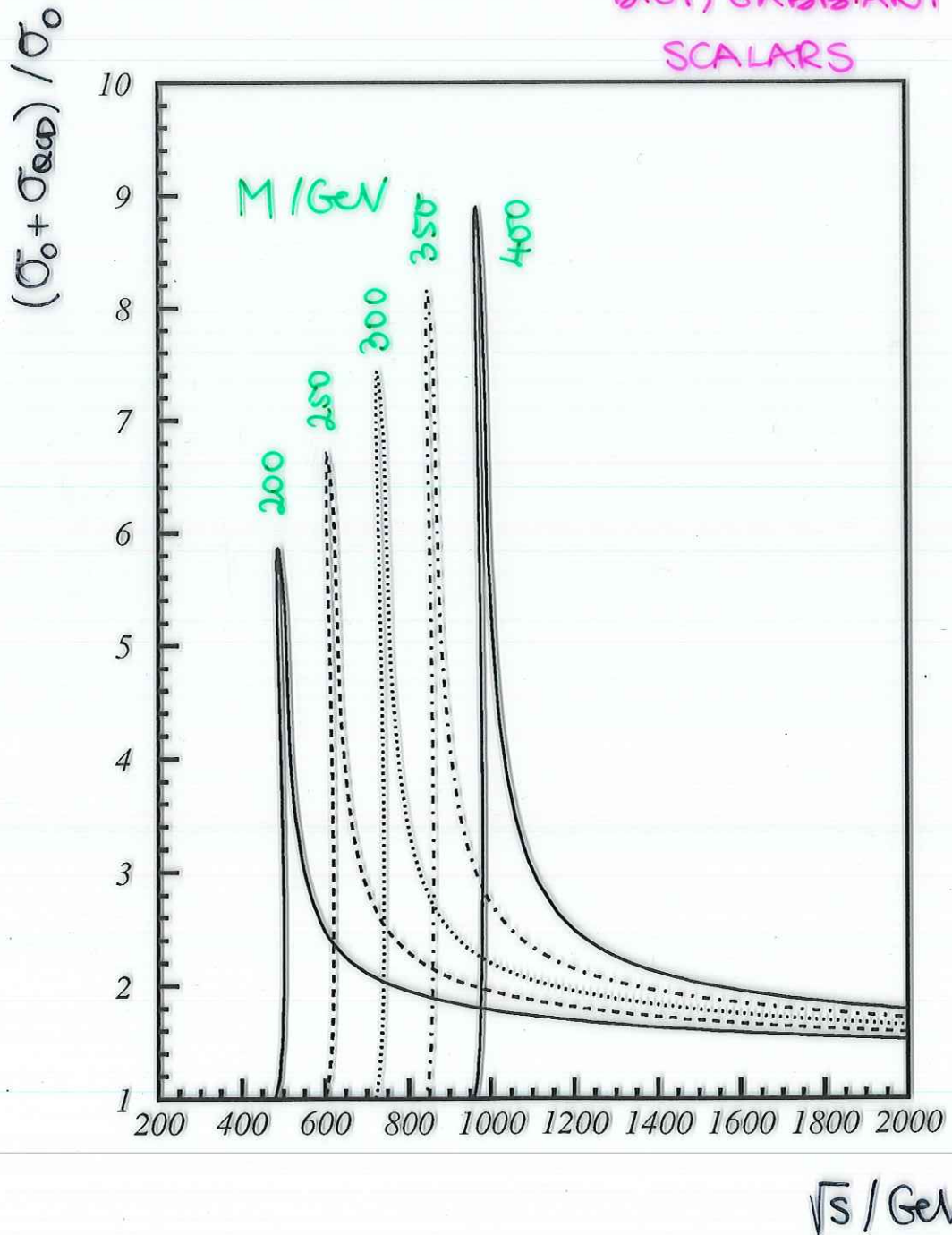
$$E = W - 2M_S, \quad \Gamma_S \simeq 1 \text{ GeV}, \quad \alpha_s = \alpha_s(M_S^2) \text{ fixed}$$

(hadronic width of S's)

3. QCD CORRECTIONS

(LIT: BIGI, FADIN, KHOZE)

BIGI, GABBIANI '91
SCALARS



↳ LEPTOQUARKONIUM ?

(JB '93)

Formation of Bound States ?

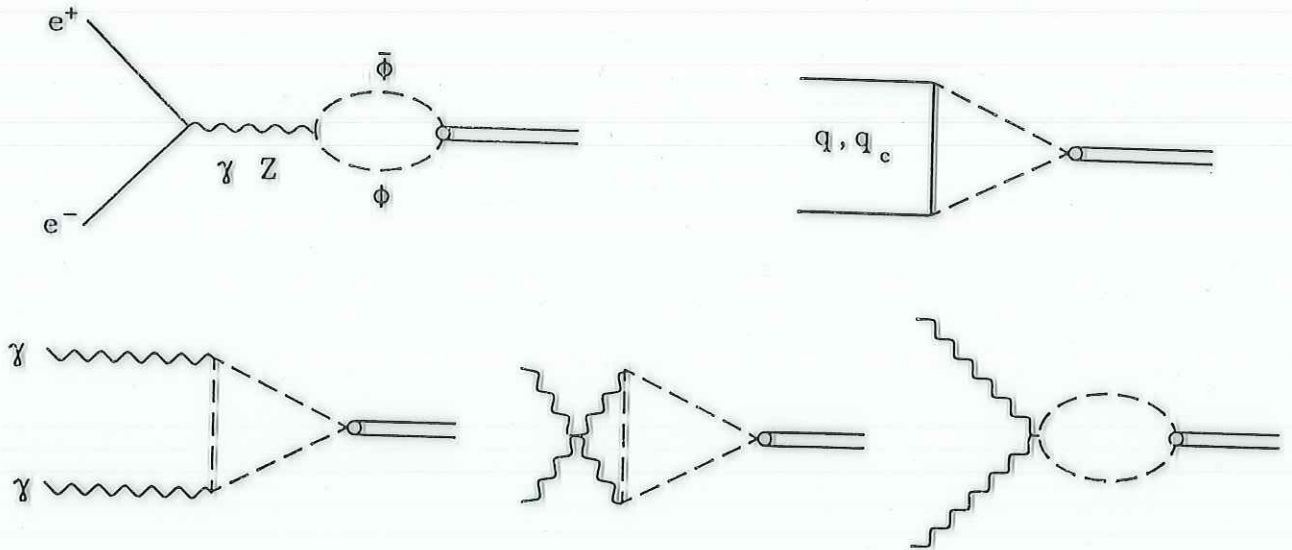
'Leptoquarkonia'

JB '93

BOWSER - CHAO et al '98

KISELEV '98

$$\beta \ll 1$$



$$\sigma(e^+e^- \rightarrow \Phi_s \bar{\Phi}_s) \propto \beta^3$$

$$\sigma(e^+e^- \rightarrow \Phi_v \bar{\Phi}_v) \propto \beta^3$$

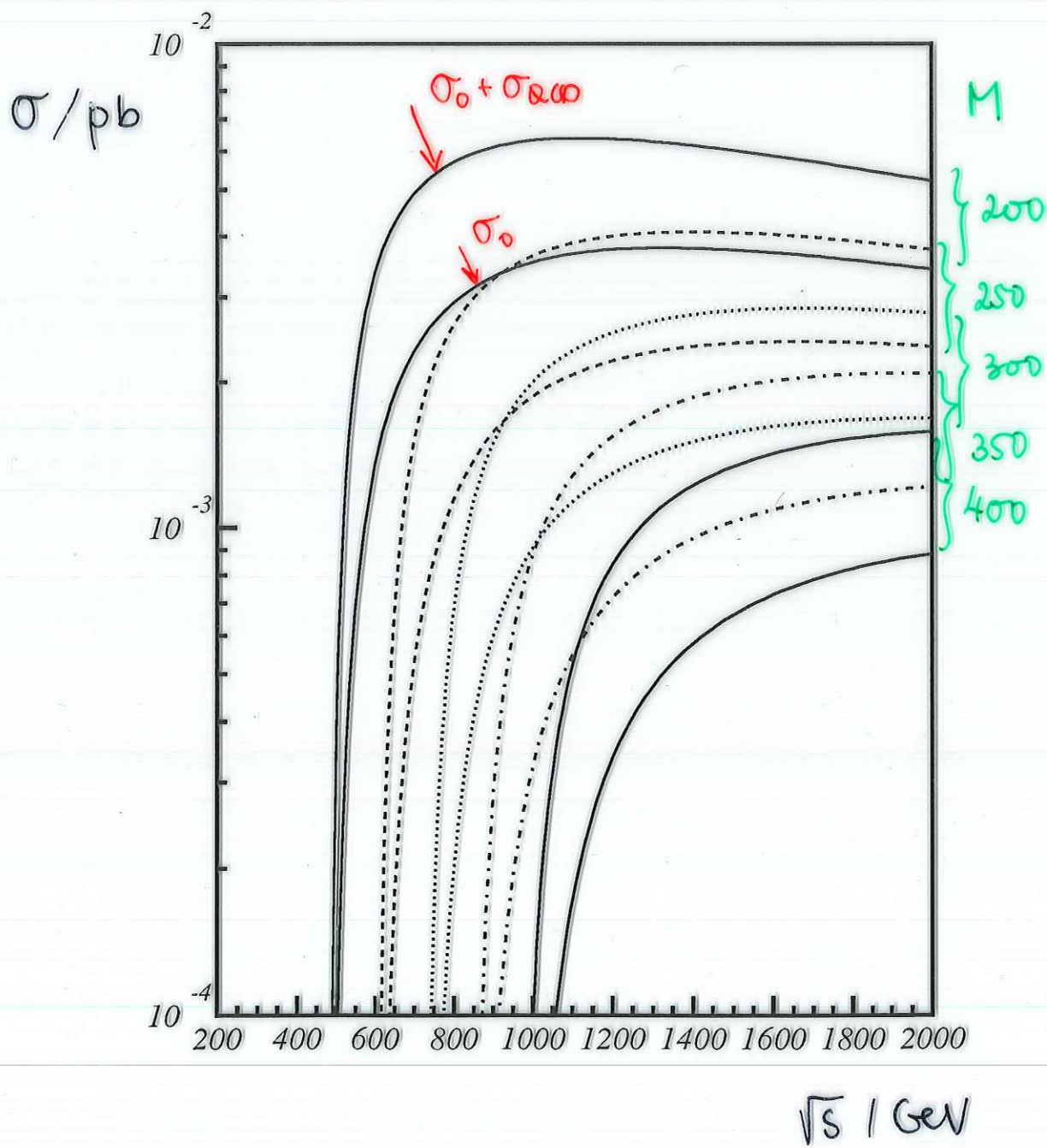
$$\sigma(\gamma\gamma \rightarrow \Phi_s \bar{\Phi}_s) \propto \beta$$

$$\sigma(\gamma\gamma \rightarrow \Phi_v \bar{\Phi}_v) \propto \beta$$

} preferred!

(14)

$$\Gamma_\phi \sim \lambda_{eq}^2 M_\phi \lesssim 100 \text{ MeV}$$



- IMPORTANT THRESHOLD FACTOR!
- COULOMB-SINGULARITY

Conclusions

1. $\sigma_{S,V}(\gamma\gamma \rightarrow \Phi\bar{\Phi}) \propto Q_{\Phi}^4$
2. Background free window to measure κ_A, λ_A iff vector LQ's exist (then found by LHC before)
3. Strong threshold enhancement $O(5)$
4. Lepoquarkonia iff $\Gamma_{\Phi} \sim 100 \text{ MeV}$
5. Final state corrections are mandatory.